# CM-07 Solved Problem* <br> Small Oscillations 

Three spring problem

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Find normal frequencies, normal modes and normal coordinates of small oscillations for three identical springs system shown in Fig 1 assuming that the system oscillates in a horizontal $X-Y$ plane. The natural length of each spring is $\sqrt{ } 2$.


Fig. 1 Three Springs

## © Solution:

- Choose generalized coordinates

The obvious generalized coordinates to be used are the Cartesian coordinates $x, y$ of the body tied to the springs.

- Kinetic and energy

Now we write the kinetic and potential energies and express them in terms of the generalized coordinates $\theta, \phi$.

$$
\begin{equation*}
\text { K.E. }=\frac{1}{2} m\left(\dot{x}^{2}+\dot{y}^{2}\right) \tag{1}
\end{equation*}
$$

[^0]- Potential energy When the body is at position $(x, y)$, the potential energy of a spring is $\frac{1}{2} k(\Delta L)^{2}$ where $\Delta L$ is the extension or compression of the spring. Therefore we get the potential energies of the three springs as

$$
\begin{align*}
& V_{1}=\frac{k}{2}\left(\sqrt{(x-1)^{2}+(y-1)^{2}}-\sqrt{ } 2\right)^{2}  \tag{2}\\
& V_{2}=\frac{k}{2}\left(\sqrt{(x+1)^{2}+(y-1)^{2}}-\sqrt{ } 2\right)^{2}  \tag{3}\\
& V_{3}=\frac{k}{2}\left(\sqrt{(x+1)^{2}+(y+1)^{2}}-\sqrt{ } 2\right)^{2} \tag{4}
\end{align*}
$$

Hence

$$
\begin{equation*}
\text { P.E. }=V_{1}+V_{2}+V_{3} \tag{5}
\end{equation*}
$$

- Write the Lagrangian The Lagrangian for the three spring problem becomes given by

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2} m \dot{x}^{2}+\frac{1}{2} m \dot{y}^{2}-V_{1}-V_{2}-V_{3} \tag{6}
\end{equation*}
$$

- Find the equilibrium points and expand the Lagrangian in powers of displacements from equilibrium. The equilibrium position of the body is obviously the origin, $x=0, y=0$. Therefore, we need to expand the potential in powers of $x, y$ and retain terms of $u p$ to only second order in $x$ and $y$. For this purpose we will use the known binomial expansion given by

$$
\begin{equation*}
(1+z)^{\alpha} \approx 1+\alpha z+\frac{\alpha(\alpha-1)}{2!} z^{2}+\frac{\alpha(\alpha-1)(\alpha-2)}{3!} z^{3}+\cdots \tag{7}
\end{equation*}
$$

- Expansion of potential energy terms $V_{1}, V_{2}, V_{3}$ in powers of $x, y$.

$$
\begin{align*}
V_{1}(x, y) & =\frac{1}{2} k\left(\sqrt{(x-1)^{2}+(y-1)^{2}}-\sqrt{ } 2\right)^{2} \\
& =\frac{1}{2} k\left\{x^{2}-2 x+1+y^{2}-2 y+1+2-2 \sqrt{ } 2 \sqrt{(x-1)^{2}+(y-1)^{2}}\right\}  \tag{8}\\
& =\frac{1}{2} k\left\{x^{2}+y^{2}+2 x+2 y+4-2 \sqrt{ } 2\left(2-2 x-2 y+x^{2}+y^{2}\right)^{1 / 2}\right\}  \tag{9}\\
& =\frac{1}{2} k\left\{x^{2}+y^{2}-2 x-2 y+4-4\left(1-x-y+\frac{1}{2}\left(x^{2}+y^{2}\right)\right)^{1 / 2}\right\} \tag{10}
\end{align*}
$$

Do a binomial expansion of $\left(1+x+y+\frac{1}{2}\left(x^{2}+y^{2}\right)\right)^{1 / 2}$ and verify that

$$
\begin{equation*}
V_{1}(x, y) \approx \frac{1}{4} k\left(x^{2}+y^{2}\right)+\frac{1}{2} k x y \tag{11}
\end{equation*}
$$

In a similar fashion, we would get

$$
\begin{align*}
& V_{2} \approx \frac{1}{4} k\left(x^{2}+y^{2}\right)+\frac{1}{2} k x y  \tag{12}\\
& V_{3} \approx \frac{1}{4} k\left(x^{2}+y^{2}\right)-\frac{1}{2} k x y \tag{13}
\end{align*}
$$

Verify These equations

- Lagrangian for small oscillations

Therefore the Lagrangian for small oscillations takes the form

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2} m \dot{x}^{2}+\frac{1}{2} m \dot{y}^{2}-\frac{3}{4}\left(x^{2}+y^{2}\right)-\frac{1}{2} k x y+\cdots . \tag{14}
\end{equation*}
$$

- We now write the equations of motion for $x$ and $y$

$$
\begin{equation*}
\frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \dot{x}}\right)=\frac{\partial \mathcal{L}}{\partial x}, \quad \frac{d}{d t}\left(\frac{\partial \mathcal{L}}{\partial \dot{y}}\right)=\frac{\partial \mathcal{L}}{\partial x} . \tag{15}
\end{equation*}
$$

The equation of motion take the form

$$
\begin{align*}
m \ddot{x} & =-k\left\{\frac{3}{2} x+\frac{1}{2} y\right\}  \tag{16}\\
m \ddot{y} & =-k\left\{\frac{1}{2} x+\frac{3}{2} y\right\} \tag{17}
\end{align*}
$$

Let us use the notation Writing EOM and proceeding this way is simpler for a system with two degrees of freedom. For systems, with more than two degrees of freedom, one should proceed differently.

- It is now helpful to write the EOM in matrix form:

$$
\binom{\ddot{x}}{\ddot{y}}=-\nu^{2}\left(\begin{array}{cc}
\frac{3}{2} & \frac{1}{2}  \tag{19}\\
\frac{1}{2} & \frac{3}{2}
\end{array}\right)\binom{x}{y}
$$

- In a normal mode of vibration all the coordinates oscillate with the same frequency. Therefore we write

$$
\begin{equation*}
x=A e^{i \omega t}, \quad y=B e^{i \omega t} . \tag{20}
\end{equation*}
$$

or in matrix form we $\binom{x}{y}=\mathrm{e}^{i \omega t}\binom{A}{B}$ Substituting (21) in (19), and canceling $e^{i \omega t}$, we get

$$
-\omega^{2}\binom{A}{B}=-\nu^{2}\left(\begin{array}{ll}
\frac{3}{2} & \frac{1}{2}  \tag{21}\\
\frac{1}{2} & \frac{3}{2}
\end{array}\right)
$$

Rearranging this equation gives

$$
\left(\begin{array}{cc}
\omega^{2}-\frac{3}{2} \nu^{2} & -\frac{1}{2} \nu^{2}  \tag{22}\\
-\frac{1}{2} \nu^{2} & \omega^{2}-\frac{3}{2} \nu^{2}
\end{array}\right)\binom{A}{B}=0 .
$$

This equation will have nontrivial solution for $A, B$ only if the determinant of the $2 \times 2$ matrix on the left is zero.

$$
\operatorname{det}\left(\begin{array}{cc}
\omega^{2}-\frac{3}{2} \nu^{2} & -\frac{1}{2} \nu^{2}  \tag{23}\\
-\frac{1}{2} \nu^{2} & \omega^{2}-\frac{3}{2} \nu^{2}
\end{array}\right)=0
$$

This determines the frequencies of the normal modes of vibration as solutions of the equation

$$
\begin{align*}
\left(\omega^{2}-\frac{3}{2} \nu^{2}\right)^{2}-\frac{1}{2} \nu^{4} & =0  \tag{24}\\
\left(\omega^{2}-2 \nu^{2}\right)\left(\omega^{2}-\nu^{2}\right) & =0 \tag{25}
\end{align*}
$$

Hence the two frequencies are given by

$$
\begin{equation*}
\omega_{1}=\nu=\sqrt{\frac{k}{m}} \quad \omega_{2}=\sqrt{ } 2 \nu=\sqrt{\frac{2 k}{m}} \tag{26}
\end{equation*}
$$

- Normal Coordinates

Now solve the equations (21), or (22), for $A, B$ for the two frequencies. For the two frequencies, this gives solutions. Write your answers for $A, B$ as column vectors

$$
\begin{array}{ll}
\omega=\omega_{1}, & \chi_{1}=N_{1}\binom{1}{-1} \\
\omega=\omega_{2}, & \chi_{2}=N_{2}\binom{1}{1}
\end{array}
$$

where $N_{1}, N_{2}$ are some normalization constants.

- We fix normalization constants using

$$
\begin{align*}
\chi_{1}^{(T)} M \chi_{1} & =1 \Longrightarrow 2 m N_{1}^{2}=1  \tag{27}\\
\chi_{2}^{(T)} M \chi_{2} & =1 \Longrightarrow 2 m N_{2}^{2}=1 \tag{28}
\end{align*}
$$

- Next we define a matrix $S$ as

$$
S=\left(\begin{array}{ll}
\chi_{1} & \chi_{2}
\end{array}\right)=\sqrt{\frac{1}{2 m_{1}}}\left(\begin{array}{cc}
1 & 1 \\
-1 & 1
\end{array}\right)
$$

- Write the coordinates $(x, y)$ in terms of normal coordinates $\left(Q_{1}, Q_{2}\right)$ and also the inverse relation

$$
\begin{equation*}
\binom{x}{y}=S\binom{Q_{1}}{Q_{2}} ; \quad\binom{Q_{1}}{Q_{2}}=S^{-1} \cdot\binom{x}{y} \tag{30}
\end{equation*}
$$

- Verify that the Lagrangian written in terms of normal coordinates becomes

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2}\left(\dot{Q}_{1}^{2}+\dot{Q}_{2}^{2}\right)-\frac{1}{2}\left(\omega_{1}^{2} Q_{1}^{2}+\omega_{2}^{2} Q_{2}^{2}\right) . \tag{31}
\end{equation*}
$$

In each normal mode of vibration only one normal coordinate varies harmonically with time, all other normal coordinates remain constant.

## Time variation of coordinates

From Eq.(31), Euler Lagrange equations for the normal coordinates are easy to write and we get

$$
\begin{equation*}
\ddot{Q}_{1}+\omega_{1}^{2} Q_{1}=0, \quad \ddot{Q}_{2}+\omega_{2}^{2} Q_{2}=0, \tag{32}
\end{equation*}
$$

and the solutions are

$$
\begin{equation*}
Q_{1}(t)=Q_{10} e^{i \omega_{1} t}+\text { c.c., } \quad Q_{2}(t)=Q_{20} e^{i \omega_{1} t}+\text { c.c. } \tag{33}
\end{equation*}
$$

Here the coordinate amplitudes $Q_{10}, Q_{20}$ are complex numbers and c.c. means complex conjugate of the first term. The time variation of the coordinates $(x, y)$ can now be written down as

$$
\begin{equation*}
x(t)=\sqrt{\frac{1}{2 m}}\left(Q_{1}(t)+Q_{2}(t)\right), \quad y(t)=\sqrt{\frac{1}{2 m}}\left(Q_{1}(t)-Q_{2}(t)\right) \tag{34}
\end{equation*}
$$

Taking $Q_{10}=A_{1}+i B_{1}, Q_{20}=A_{2}+i B_{2}$ The four (real) unknown constants $A_{1}, A_{2}, B_{1}, B_{2}$, can be determined if initial conditions on position and velocity vectors, $(x, y)$ and $(\dot{x}, \dot{y})$, are given.
$\otimes_{0}$ In general the time variation of Cartesian coordinates $x, y$ is superposition of both the frequencies. The motion of the body will trace out a Lissajous figure.
$\otimes_{0}$ Only with special initial conditions, the coordinates will have harmonic time variation with a single frequency.

## Questions for you

- Write explicit expressions for most general solution, in terms of the four unknown constants $A_{1}, B_{1}, A_{2}, B_{2}$, for the time dependence of normal coordinates. Do the same for the coordinates $(x, y)$.
- Write initial conditions so the the system may vibrate in one of the normal modes. These modes of vibration are called modes of vibration.
- Geometrically describe the motion of the body in each of the two normal modes.
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