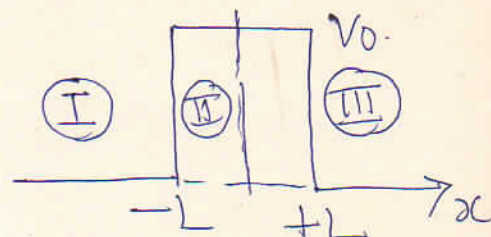


Computational Details.

Consider a beam of particles incident on a target which may be represented by a potential $V(x)$

$$V(x) = \begin{cases} 0 & x < -L \\ V_0 & -L < x < L \\ 0 & x > L \end{cases}$$



for positive V_0 , we call such a potential as square barrier.

Classically we expect that

$E < V_0$ the particles will get reflected.

$E > V_0$ the particles will get transmitted.

We want to solve the Schrodinger equation and obtain energy eigenfunctions. We shall write the solutions in the three regions

$$\textcircled{\text{I}} : x < -L \quad \textcircled{\text{II}} : -L < x < L \quad \textcircled{\text{III}} : x > L$$

separately. The solutions are given by, for $E < V_0$

$$u_{\text{I}}(x) = A e^{ikx} + B e^{-ikx}$$

$$u_{\text{II}}(x) = C e^{\beta x} + D e^{-\beta x}, \quad \beta^2 = \frac{2m(V_0 - E)}{\hbar^2}$$

$$u_{\text{III}}(x) = F e^{ikx} + G e^{-ikx}, \quad k^2 = \frac{2mE}{\hbar^2}$$

We the incident beam is coming from $-\infty$, there will be no particles travelling to the left in region III. So we set $G = 0$.

Impose the boundary conditions at $x=L$ that the wave function and its derivative must be continuous.

This means

$$u_{II}(L) = u_{III}(L) \quad ; \quad u'_{II}(L) = u'_{III}(L)$$

$$\text{or } ce^{\beta L} + De^{-\beta L} = Fe^{ikL} + Ge^{-ikL} \quad \text{--- ①}$$

$$\beta ce^{\beta L} - \beta De^{-\beta L} = Fike^{ikL} - ikGe^{-ikL} \quad \text{--- ②}$$

Set $G=0$; multiply ① by ik and subtract from ②

to get

$$(\beta - ik)ce^{\beta L} - (\beta + ik)De^{-\beta L} = 0$$

$$\text{or } \frac{D}{c} = \left(\frac{\beta - ik}{\beta + ik} \right) e^{2\beta L} \quad \text{--- ③}$$

--- x ---

Boundary conditions at $x=-L$ are

$$u_I(L) = u_{II}(L) \quad ; \quad u'_I(L) = u'_{II}(L) \quad \text{--- ④}$$

These equations give

$$ce^{-\beta L} + De^{\beta L} = Ae^{-ikL} + Be^{ikL} \quad \text{--- ⑤}$$

$$c\beta e^{-\beta L} - D\beta e^{\beta L} = Aike^{-ikL} - Blike^{ikL} \quad \text{--- ⑥}$$

We will solve these equations for B/A . Multiply ⑤ by ik and add ⑥ to get

$$c(\beta + ik)e^{-\beta L} - D(\beta - ik)e^{\beta L} = 2Aike^{-ikL} \quad \text{⑦}$$

Next multiply ⑤ by ik and subtract from ⑥

$$c(\beta - ik)e^{-\beta L} - D(\beta + ik)e^{\beta L} = -2ikBe^{ikL} \quad \text{⑧}$$

Use (3) to eliminate D from (7). This gives

$$\begin{aligned}
2Aik e^{-ikL} &= C(\beta+ik)e^{-\beta L} - \frac{(\beta-ik)^2}{(\beta+ik)} e^{3\beta L} C \\
&= C \frac{(\beta+ik)^2 e^{-\beta L} - (\beta-ik)^2 e^{3\beta L}}{(\beta+ik)} \\
&= \frac{C e^{\beta L}}{(\beta+ik)} \left[(\beta+ik) e^{-2\beta L} - (\beta-ik) e^{2\beta L} \right] \\
&= \frac{C e^{\beta L}}{(\beta+ik)} \left((\beta^2 - k^2) (e^{-2\beta L} - e^{2\beta L}) + 2i\beta k (e^{-2\beta L} + e^{2\beta L}) \right) \\
&= \frac{C e^{\beta L}}{(\beta+ik)} \left(-2(\beta^2 - k^2) \sinh 2\beta L + 4i\beta k \cosh 2\beta L \right) \quad \text{--- (9)}
\end{aligned}$$

Next use (3) again, this time to eliminate D from (8). This gives

$$\begin{aligned}
2ikB e^{ikL} &= D(\beta+ik)e^{\beta L} - C(\beta-ik)e^{-\beta L} \\
&= \left(\frac{\beta-ik}{\beta+ik} \right) C(\beta+ik) e^{3\beta L} - C(\beta-ik) e^{-\beta L} \\
&= C(\beta-ik) e^{\beta L} (e^{2\beta L} - e^{-2\beta L}) \\
&= C(\beta-ik) e^{\beta L} 2 \sinh 2\beta L \quad \text{--- (10)}
\end{aligned}$$

Divide (9) (10) by (9) to get

$$e^{2ikL} \left(\frac{B}{A} \right) = \frac{2(\beta-ik)(\beta+ik) \sinh 2\beta L}{-2(\beta^2 - k^2) \sinh 2\beta L + 4i\beta k \cosh 2\beta L}$$

$$\text{or } \left(\frac{B}{A} \right) = \frac{(\beta^2 + k^2) \sinh 2\beta L e^{-2ikL}}{(k^2 - \beta^2) \sinh 2\beta L + 2i\beta k \cosh 2\beta L} \quad \text{--- (11)}$$

Consider Eq (1) and (5) and solve for (C/A) and (F/C)

Eq (1) with $G=0$

$$c e^{\beta L} + D e^{-\beta L} = F e^{i k L} \quad (12)$$

Eq (5)

$$c e^{-\beta L} + D e^{\beta L} = A e^{-i k L} + B e^{i k L} \quad (13)$$

(12) gives

$$F e^{i k L} = c e^{\beta L} + D e^{-\beta L}$$

$$= c e^{\beta L} + \left(\frac{\beta - i k}{\beta + i k} \right) c e^{\beta L}$$

|| use (3) and eliminate D

$$F = c e^{\beta L} \frac{2\beta}{(\beta + i k)} e^{-i k L} \quad (14)$$

Next solve for c/A:

L.H.S of (13)

$$= c e^{-\beta L} + D e^{\beta L}$$

$$= c e^{-\beta L} + c \left(\frac{\beta - i k}{\beta + i k} \right) e^{3\beta L}$$

|| use (3) and eliminate D

$$= \frac{c e^{\beta L}}{(\beta + i k)} \left((\beta + i k) e^{-2\beta L} + (\beta - i k) e^{2\beta L} \right)$$

$$= \frac{c e^{\beta L}}{(\beta + i k)} \left(\beta (e^{-2\beta L} + e^{2\beta L}) - i k (e^{2\beta L} - e^{-2\beta L}) \right)$$

$$= \frac{c e^{\beta L}}{(\beta + i k)} \left(2\beta \cosh 2\beta L - 2i k \sinh 2\beta L \right) \quad (15)$$

R.H.S. of (13) gives

$$A e^{-i k L} + B e^{i k L} =$$

Substitute for B from (11)

$$= A e^{-i k L} + \frac{A e^{-i k L} (\beta + k^2) \sinh 2\beta L}{(k^2 - \beta^2) \sinh 2\beta L + 2i \beta k \cosh 2\beta L}$$

$$= A e^{-ikL} \left(1 + \frac{(\beta^2 + k^2) \sinh 2\beta L}{(k^2 - \beta^2) \sinh 2\beta L + 2i\beta k \cosh 2\beta L} \right)$$

$$= A e^{-ikL} \frac{2k^2 \sinh 2\beta L + 2i\beta k \cosh 2\beta L}{(k^2 - \beta^2) \sinh 2\beta L + 2i\beta k \cosh 2\beta L} \quad (16)$$

Equating L.H.S. of (13) and R.H.S. of (13) using expressions (15) and (16) gives

$$\frac{C e^{\beta L}}{(\beta + ik)} (2\beta \cosh 2\beta L - ik \sinh 2\beta L) \\ = A e^{-ikL} \frac{(ik) (2\beta \cosh 2\beta L - 2ik \sinh 2\beta L)}{(k^2 - \beta^2) \sinh 2\beta L + 2i\beta k \cosh 2\beta L}$$

$$\text{or } \frac{C}{A} = (\beta + ik) e^{-(\beta + ik)L} \times \frac{ik}{(k^2 - \beta^2) \sinh 2\beta L + 2i\beta k \cosh 2\beta L} \quad (17)$$

$$\text{From (14)} \\ \frac{F}{C} = e^{(\beta - ik)L} \frac{2\beta}{(\beta + ik)} \quad \dots \quad (18)$$

multiplying (17) and (18) gives.

$$\frac{F}{A} = \frac{2i\beta k e^{-2ikL}}{(k^2 - \beta^2) \sinh(2\beta L) + 2i\beta k \cosh(2\beta L)} \quad (19)$$

Transmission coefficient is $|F/A|^2$ and we get

$$T = |F/A|^2 = \left(1 + \frac{(\beta^2 + k^2)^2}{2\beta^2 k^2} \sinh^2 2\beta L \right)^{-1} \quad (20)$$

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Square barrier - - continued.

$E < V_0$: Very high barrier

Transmission Coefficient -

$$= \left(1 + \frac{(\beta^2 + k^2)^2}{4\beta^2 k^2} \sinh^2 2\beta L \right)^{-1}$$

$$\beta^2 = \frac{2m(V_0 - E)}{\hbar^2}$$

$$k^2 = \frac{2mE}{\hbar^2}$$

$E < V_0$

$$\beta^2 + k^2 = \frac{2mV_0}{\hbar^2}$$

$$2\beta^2 k^2 = \frac{8(V_0 - E)E m^2}{\hbar^4}$$

$$T = \left(1 + \frac{V_0^2}{4(V_0 - E)E} \frac{e^{+4\beta L}}{4} \right)^{-2}$$

$$\begin{aligned} & \sinh 2\beta L \\ &= \frac{1}{2} e^{2\beta L} \\ &\approx \frac{1}{2} \end{aligned}$$

$$\approx \frac{16E(V_0 - E)}{V_0^2} e^{-4\beta L}$$

$E \approx V_0$ $E \approx V_0$ $\beta L \approx 0$

$$T = \left(1 + \frac{V_0^2}{4(V_0 - E)E} 4\beta^2 L^2 \right)^{-1} \quad \sinh 2\beta L \approx 2\beta L$$

$$= \left(1 + \frac{V_0}{4(V_0 - E)E} \times 4 \frac{2m(V_0 - E)L^2}{\hbar^2} \right)^{-1}$$

$$= \left(1 + \frac{V_0^2 L^2 2m}{E\hbar^2} \right)^{-1} = \left(1 + \frac{V_0^2}{E\Delta} \right)^{-1}$$

$$\Delta = \frac{\hbar^2}{2mL^2}$$

$E > V_0$

$$\beta^2 = \frac{2m(V_0 - E)}{\hbar^2} = -\alpha^2$$

$$\alpha^2 = \frac{2m(E - V_0)}{\hbar^2} \quad \beta \rightarrow i\alpha$$

$$\begin{aligned} \sinh 2\beta L &= \frac{e^{2\beta L} - e^{-2\beta L}}{2} \\ &= \frac{e^{2i\alpha L} - e^{-2i\alpha L}}{2} \end{aligned}$$

$$= i \sin 2\alpha L$$

$$\sinh^2 2\beta L = -\sin^2 2\alpha L$$

$$\therefore T = \left(1 + \frac{V_0^2}{4(V_0 - E)E} \times (-\sin^2 2\alpha L) \right)^{-1}$$

$$= \left(1 + \frac{V_0^2}{4(E - V_0)E} \sin^2 2\alpha L \right)^{-1}$$

maximum when $2\alpha L = n\pi$

$$\alpha^2 = \frac{2m(E - V_0)}{\hbar^2} = \frac{p^2}{\hbar^2} = \frac{4\pi^2}{\lambda^2}$$

$$\left(\frac{2L}{\lambda} \right) = n\pi \quad \alpha = \frac{2\pi}{\lambda}$$

$$(2L) \left(\frac{2\pi}{\lambda} \right) = n\pi$$

$$(2L) = n \left(\frac{\lambda}{2} \right)$$

Integral multiples of $\lambda/2$ fit into the well range.
then $T = 1$.

$$E < V_0$$

$$T = \left(1 + \frac{\sinh^2 \beta L}{E/V_0 (1 - E/V_0)} \right)^{-1}$$

When $E < V_0$ T increases monotonically with E .

$$\beta L \gg 1 \quad \sinh^2 \beta L \approx \frac{1}{4} e^{\beta L} \gg 1$$

$$\begin{aligned} T &= 16 \frac{E}{V_0} \left(1 - \frac{E}{V_0} \right) e^{-\beta L} \\ &= 16 \left(\frac{E}{V_0} \right) \left(1 - \frac{E}{V_0} \right) e^{-4 \sqrt{\frac{V_0 - E}{\Delta}}} \end{aligned}$$

$$\Delta = \frac{\hbar^2}{2mL^2} \quad \text{uncertainty energy for particle confined to length } L.$$

$$\beta L \gg 1 \quad V_0 - E \gg \Delta$$

$$E > V_0 \quad \beta \rightarrow q$$

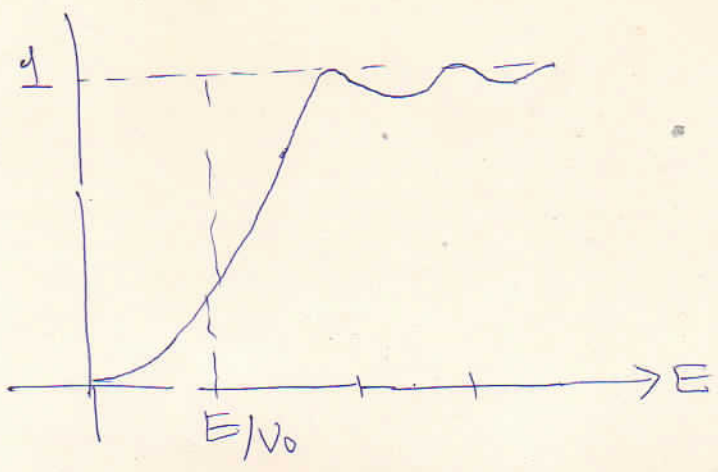
$$T(E) = \left(1 + \frac{\sin^2 \alpha L}{4 E/V_0 (1 - E/V_0)} \right)^{-1}$$

$$\alpha^2 = \frac{2m}{\hbar^2} (E - V_0) L$$

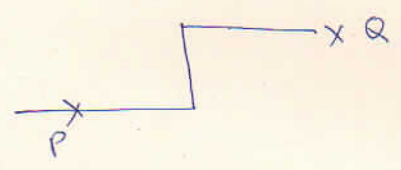
$$T(E) = 1 \quad \text{when} \quad \sin \alpha L = 0$$

$$\alpha L = n\pi$$

$$\Rightarrow \frac{2\pi}{\lambda} \cdot L = n\pi, \quad \lambda = n \cdot \lambda/2$$



Definition of transmission and reflection coefficient when potential does not go to zero at infinity.



No of particles received at a point P. per sec

$$= N$$

$$N e^{iKx}$$

$$= v_1 \rho = (N)^2 v_1$$

$$T = \left| \frac{N_2}{N_1} \right|^2 \left(\frac{v_2}{v_1} \right)$$