

Notes for Lectures in Quantum Mechanics *

Integral Equation for Scattering

A. K. Kapoor
<http://0space.org/users/kapoor>
akkapoor@cmi.ac.in; akkhcu@gmail.com

Contents

1	Green function for free particle Schrodinger Equation	1
2	Set up integral equation	1
3	Verify boundary condition	2
4	Scattering amplitude	3

1 Green function for free particle Schrodinger Equation

In order to convert the Schrodinger equation

$$\left\{ \frac{-\hbar^2}{2\mu} \nabla^2 + v(r) \right\} \psi = E\psi \quad (1)$$

into an integral equation we first rewrite it as

$$(\nabla^2 + k^2) \psi = u(\vec{r})\psi(\vec{r}) \quad (2)$$

where $k^2 = \frac{2\mu E}{\hbar^2}$, $U(r) = \frac{2\mu}{\hbar^2}V(r)$ and we have defined Green function $G(\vec{r})$ as a solution of

$$(\nabla^2 + k^2) G(\vec{r}) = -\delta(\vec{r}). \quad (3)$$

Two possible solutions for Green function are $G(\vec{r})$ are

$$G_{\pm}(\vec{r}) = \frac{e^{\pm ikr}}{4\pi r}. \quad (4)$$

Using the Green function we can now write down a "formal" solution of Eq.(2) as a solution of integral equation.

2 Set up integral equation

Using a Green function which is a solution of Eq.(3) a formal solution for Eq.(2) can be written as

$$\psi(\vec{r}) = \phi(\vec{r}) - \int G(|\vec{r} - \vec{r}'|)U(|\vec{r}'|)\psi(\vec{r}')d^3r' \quad (5)$$

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where $\phi(\vec{r})$ is a solution of the equation

$$(\nabla^2 + k^2) \phi(\vec{r}) = 0. \quad (6)$$

It turns out that the choice

$$G(\vec{r}) = \frac{e^{ikr}}{4\pi r} \quad (7)$$

for the Green function leads to the correct scattering boundary condition on the wave function.

3 Verify boundary condition

For the scattering problem we must select the function $\phi(\vec{r})$ to be

$$\phi(\vec{r}) = \exp(i\vec{k}_i \cdot \vec{r}) \quad (8)$$

where \vec{k}_i is the momentum of the incident particles. Substituting Eq.(7) Eq.(8) in Eq.(5) we get the integral equation for the scattering to be

$$\psi(\vec{r}) = \exp(i\vec{k}_i \cdot \vec{r}) - \frac{1}{4\pi} \int \frac{e^{ik|\vec{r}-\vec{r}'|}}{|\vec{r}-\vec{r}'|} U(\vec{r}') \psi(\vec{r}') d^3 r'. \quad (9)$$

To verify that $\psi(\vec{r})$ given by Eq.(9) does indeed have correct asymptotic property, we need to expand $|\vec{r}-\vec{r}'|$ in powers of $\frac{\vec{r}}{r'}$ we shall assume that the potential is short range potential so that the contribution to integral over \vec{r}' comes from small value of r' .

Some algebraic manipulations

Expand $|\vec{r}-\vec{r}'|$ in powers of r'

$$|\vec{r}-\vec{r}'| = \sqrt{r^2 + r'^2 - 2\vec{r} \cdot \vec{r}'} \quad (10)$$

$$= r \left(1 - 2\frac{\vec{r} \cdot \vec{r}'}{r^2} + \frac{r'^2}{r^2} \right)^{1/2}. \quad (11)$$

Using binomial expansion we get

$$|\vec{r}-\vec{r}'| = r \left(1 - \frac{\vec{r} \cdot \vec{r}'}{r^2} + O\left(\frac{r'^2}{r^2}\right)^2 \right). \quad (12)$$

Large r expansion of the formal solution (9)

We substitute Eq.(12) in the exponential and in the factor $\frac{1}{|\vec{r}-\vec{r}'|}$ in Eq.(9) and

write $\frac{1}{|\vec{r}-\vec{r}'|} \approx \frac{1}{r}$ to get

$$\psi(\vec{r}) \longrightarrow \exp(i\vec{k}_i \cdot \vec{r}) - \frac{1}{4\pi r} \int \exp\left(ikr - ik\frac{\vec{r} \cdot \vec{r}'}{r^2}\right) U(\vec{r}') \psi(\vec{r}') d^3 r' \quad (13)$$

$$= \exp(i\vec{k}_i \cdot \vec{r}) - \frac{e^{ikr}}{4\pi r} \int \exp(-ik\hat{n} \cdot \vec{r}') U(\vec{r}') \psi(\vec{r}') d^3 r'. \quad (14)$$

In the last step we have introduced a unit vector $\hat{n} = \vec{r}/r$. The Eq.(14) gives the probability amplitude (wave function) at \vec{r} . If the particles are to reach at a detector at \vec{r} , the vector \hat{n} must be in the direction of the final momentum and parallel to k_f . Note that

$$|\vec{k}_i| = |\vec{k}_f| = k \quad (15)$$

holds as a consequence of energy conservation and hence

$$k(\hat{n} \cdot \vec{r}') = \vec{k}_f \cdot \vec{r}'. \quad (16)$$

Thus Eq.(14) takes the form

$$\psi(\vec{r}) \approx \exp(i\vec{k}_i \cdot \vec{r}) - \frac{e^{ikr}}{4\pi r} \int \exp(-i\vec{k} \cdot \vec{r}') U(r') \psi(r') d^3 r'. \quad (17)$$

This asymptotic behaviour is of the form expected for large r

$$\psi(\vec{r}) \approx \exp(i\vec{k}_i \cdot \vec{r}) - \frac{e^{ikr}}{r} f(\theta, \phi). \quad (18)$$

4 Scattering amplitude

Comparing Eq.(17) with Eq.(18) we see that the scattering amplitude is given by

$$f(\theta, \phi) = -\frac{1}{4\pi} \int \exp(-i\vec{k}_f \cdot \vec{r}') U(r') \psi(r') d^3 r' \quad (19)$$

$$= -\left(\frac{\mu}{2\pi\hbar^2}\right) \int \exp(-i\vec{k}_f \cdot \vec{r}') V(r') \psi(r') d^3 r'. \quad (20)$$

It must be noted that the integral equation Eq.(9) and the expression for the scattering amplitude in Eq.(20) are **exact** results.