

Question: Approaching the problem of electric field for a uniformly charged sphere in two different ways, find the value of the integral

$$\vec{\mathcal{E}} = \frac{1}{4\pi\epsilon_0} \iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3}. \quad (1)$$

⊙ **Solution:**

Method 1:

A charged sphere has radius R and centre at the origin and having spherically symmetric charge distribution ρ . The electric field for the sphere by at a point $vecr$ can be computed by using Gauss law and is given by

$$\vec{E} = \begin{cases} \frac{Q(r)}{4\pi\epsilon_0} \vec{r}, & \text{if } r > R, \\ \frac{Q(r)}{4\pi\epsilon_0} \frac{\vec{r}}{r^3}, & \text{if } r \leq R. \end{cases} \quad (2)$$

where $Q(r)$ is the total charge inside a sphere of radius and is

$$Q(r) = \int_0^r \rho(r) 4\pi r^2 dr.$$

In case of uniform charge density, $\rho(r)$ independent of r , the above result for the electric field reduces to

$$\vec{E} = \begin{cases} \frac{\rho \vec{r}}{3\epsilon_0} & \text{if } , r < R, \\ \frac{\rho \vec{r}}{3\epsilon_0} \times \frac{R^3}{r^3}, & \text{if } r < R. \end{cases} \quad (3)$$

On the other hand the Coulomb's law with superposition principle give the electric field as

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \iiint d^3r' \rho \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = \frac{\rho}{4\pi\epsilon_0} \iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} \quad (4)$$

for the case of uniformly charged sphere.

Equating the electric field answers from (3) and (4) and canceling ρ we get result

$$\frac{1}{4\pi\epsilon_0} \iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = \begin{cases} \frac{\vec{r}}{3\epsilon_0}, & \text{if } r < R, \\ \frac{\vec{r}}{3\epsilon_0} \times \frac{R^3}{r^3}, & \text{if } r < R. \end{cases} \quad (5)$$

Method 2:

This method makes use of generating function of Legendre polynomials

$$\frac{1}{(1 - 2xt + t^2)^{1/2}} = \sum_0^{\infty} t^n P_n(x). \quad (6)$$

We first note that the

$$\frac{\vec{r} - \vec{r}'}{|\vec{r} - \vec{r}'|^3} = -\nabla \frac{1}{|\vec{r} - \vec{r}'|} \quad (7)$$

Therefore required integral can be rewritten as

$$\iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = -\nabla \iiint d^3r' \frac{1}{|\vec{r} - \vec{r}'|^3} \equiv -\nabla f(r) \quad (8)$$

$$\text{where } f(r) = \iiint d^3r' \frac{1}{|\vec{r} - \vec{r}'|}. \quad (9)$$

and the integral is to be evaluated over a sphere of radius R . because the term with \vec{r}' in the numerator will vanish when averaged over all directions. Next we consider the two cases when (i) $r > R$ (ii) $r < R$ and use the generating function to expand $1/|\vec{r} - \vec{r}'|$ as follows

$$\frac{1}{|\vec{r} - \vec{r}'|} = \frac{1}{r} \sum_n \left(\frac{r'}{r}\right)^n P_n(\cos \theta), \quad \text{if } r > r' \quad (10)$$

$$= \frac{1}{r'} \sum_n \left(\frac{r}{r'}\right)^n P_n(\cos \theta) \quad \text{if } r < r' \quad (11)$$

. All terms involving Legendre polynomials $P_n(\cos \theta)$, $n \neq 0$ vanish due to orthogonality property of Legendre polynomials. Thus we get

$$f(r) = \int d\Omega \int_0^R \frac{r'^2}{r} = \frac{4\pi R^3}{3r} \quad (12)$$

Substituting the above value in (8) we get the required integral

$$\frac{1}{4\pi\epsilon_0} \iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = -\frac{1}{4\pi\epsilon_0} \nabla \left(\frac{4\pi R^3}{3r} \right) = \frac{1}{3\epsilon_0} \frac{R^3}{r^3}. \quad (13)$$

Case (ii):

In this case also only the first term survives in the expansion in Legendre polynomials. However, the radial expansion needs to be broken up as a sum

of integral over ranges $(0 < r' < r)$ and $(r < r' < R)$. Using appropriate expression from (10) and (11), we get

$$f(r) = 4\pi \left[\int_0^r (1/r) r'^2 dr' + \int_r^R (1/r') r'^2 dr' \right] \quad (14)$$

$$= 4\pi \left[\frac{r^2}{3} + \frac{R^2 - r^2}{2} \right] = (4\pi) \left(\frac{R^2}{2} - \frac{r^2}{6} \right) \quad (15)$$

Taking gradient of $f(r)$, (8) gives the final answer as

$$\frac{1}{4\pi\epsilon_0} \iiint d^3r' \frac{(\vec{r} - \vec{r}')}{|\vec{r} - \vec{r}'|^3} = \frac{\vec{r}}{3\epsilon_0}. \quad (16)$$